

Research Article

Concept of Recombination Velocity S_{fcc} at the Junction of a Bifacial Silicon Solar Cell, in Steady State, Initiating the Short-Circuit Condition

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Abstract: The aim of this study is to present technics to determine the junction recombination velocity of a bifacial polycrystalline silicon solar cell under both, constant multispectral illumination and steady short-circuit condition.

Keywords: Bifacial silicon solar cell, junction recombination velocity, short-circuit

INTRODUCTION

Our contribution consists in proposing technics that determine the junction recombination velocity S_{fcc} of a bifacial polycrystalline silicon solar cell in steady state and under constant multispectral light, initiating the short-circuit condition. Thus, we basically present a theoretical survey, in which we schematize an isolated grain of the solar cell (Blakers *et al.*, 1989; Huljić *et al.*, 2001).

As the shunt resistance R_{sh} is expressed as a calibrated function of S_f while the solar cell is under short-circuit, represented by S_{fcc} , the determination of S_{fcc} lead to the exact value of R_{sh} (Mbodji *et al.*, 2010). In solving the continuity equation, we define the expressions of the charge carriers' density (Endrös and Martinelli, 1997) and deduce the photocurrents that let us introduce the concept of the recombination velocities at the junction yielding the short-circuit photocurrent density (Mbodji *et al.*, 2012). Then we will discuss the results that we have obtained before we make a conclusion.

Theory: On behalf of the study, the solar cell placed under different levels of light is illuminated on its front, back and then simultaneously on its two sides. The model is treated as one dimensional with the origin of the solar cell junction.

During the direct conversion of the radiation energy into electric energy, two important phenomena influence the solar cell's efficiency.

These are absorption of the light (photons) by the cell and generation of electron hole pairs. These absorption and generation phenomena are of a paramount importance in determining the solar cell's electronic and electric parameters and are governed by the equation of continuity (1), and this equation lets

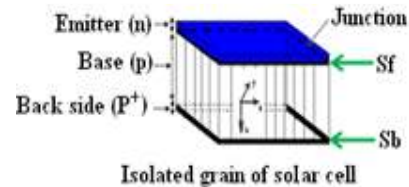


Fig. 1: Polycrystalline solar cell's structure

determine the minority carriers' density of the photo-generated charges according to the illumination mode.

In these conditions and in steady state, the equation of continuity for a one-dimensional study to which obey these charge minority carriers in the base of the x-axis is announced as follows (Fig. 1):

$$\frac{\partial^2 \delta_a(x)}{\partial x^2} - \frac{\delta_a(x)}{L^2} = -\frac{1}{D} G_a(x) \tag{1}$$

$\delta_a(x)$ is the excess minority carriers density according to the depth x in the base following the different illumination modes; α symbolizes the illumination mode with:

$\alpha = 1, \alpha = 2, \alpha = 3$ for a respective light of the front side, back side and simultaneously both sides of the solar cell.

L and D are respectively the excess minority carrier's diffusion length and coefficient in the base.

$G_a(x)$ is the generation rate of the excess minority carriers to an x distance in the base of the solar cell.

The expression of $G_a(x)$ depends on the x depth of light absorption in the base and can be written in the form of the following equation:

$$G_a(x) = n \sum_{i=1}^3 a_i \left[\xi_\alpha e^{-b_i x} + \chi_\alpha e^{-b_i(H-x)} \right] \tag{2}$$

Table1: The illumination mode parameters

Illumination mode	ξ_α	χ_α
Front side illumination $\alpha = 1$	1	0
Back side illumination $\alpha = 2$	0	1
Both sides simultaneous illumination $\alpha = 3$	1	1

H = thickness of the base;

The parameters a_i and b_i stem from the modeling of the incident illumination as defined by (Mohammad, 1987) under A.M 1.5 condition. where, $a_1 = 6.13 \times 10^{20} \text{ cm}^{-3}/\text{s}$; $a_2 = 0.54 \times 10^{20} \text{ cm}^{-3}/\text{s}$; $a_3 = 0.0991 \times 10^{20} \text{ cm}^{-3}/\text{s}$; $b_1 = 6630 \text{ cm}^{-1}$; $b_2 = 103 \text{ cm}^{-1}$; $b_3 = 130 \text{ cm}^{-1}$; n = the number of sun linking the real incident power to the reference power for a given solar spectrum. It is defined as follows:

$$n = \frac{I_{cci}}{I_{cco}} \quad (3)$$

With, I_{cci} the short-circuit photocurrent density obtained from a solar cell under A.M 1.5 condition and I_{cco} is the short-circuit current reference measured for given illumination light. The coefficients ξ_α and χ_α are defined following the solar cell's illumination modes. They are given in Table 1:

The solution $\delta_\alpha(x)$ of the carriers' diffusion Eq. (1) can be made in the following forms:

$$\delta_\alpha(x) = A_\alpha \cdot ch\left(\frac{x}{L}\right) + B_\alpha \cdot sh\left(\frac{x}{L}\right) - \sum_{i=1}^3 K_i \cdot (\xi_\alpha e^{-b_i \cdot x} + \chi_\alpha e^{-b_i \cdot (H-x)}) \quad (4)$$

The coefficients A_α and B_α ($\alpha = 1, 2, 3$) can be determined from the boundary conditions at the junction and at the back side (Bowden and Rohatgi, 2001; Möller, 1993; Alain, 1997) defined as follows:

At the junction emitter-base (for $x = 0$):

$$\frac{\partial \delta_\alpha(x)}{\partial x} \Big|_{x=0} = \frac{S_{f\alpha}}{D} \cdot \delta_\alpha(0) \quad (5)$$

At the back side of the base (for $x = H$):

$$\frac{\partial \delta_\alpha(x)}{\partial x} \Big|_{x=H} = -\frac{S_{b\alpha}}{D} \cdot \delta_\alpha(H) \quad (6)$$

$S_{f\alpha}$ and $S_{b\alpha}$ represent respectively the recombination velocities of the excess minority carriers

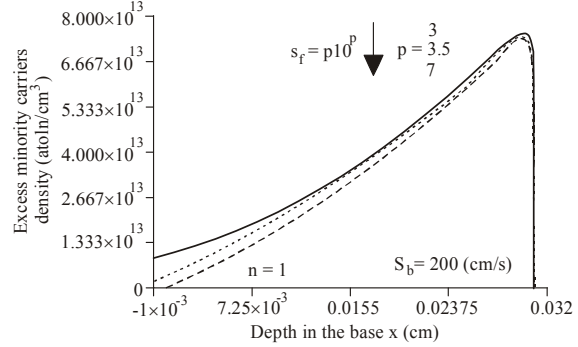


Fig. 2: Carriers' density according to the x depth in the base for a front side illumination ($S_f = p \times 10^p \text{ cm/s}$, $S_b = 200 \text{ cm/s}$, $D = 26 \text{ cm}^2/\text{s}$, $L = 0.015 \text{ cm}$ and $n = 1$)

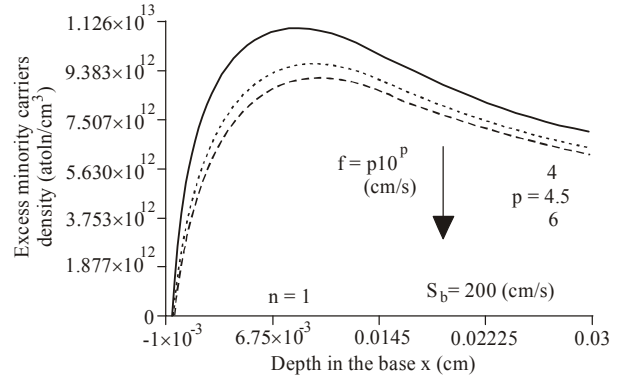


Fig. 3: Carriers' density according to the x depth in the base for a back side illumination ($S_f = p \times 10^p \text{ cm/s}$ with $S_b = 200 \text{ cm/s}$, $D = 26 \text{ cm}^2/\text{s}$, $L = 0.015 \text{ cm}$ and $n = 1$)

at the junction and at the back side (Diallo *et al.*, 2008). And K_i is defined as follows:

$$K_i = \frac{nL^2 \cdot a_i}{D(b_i^2 \cdot L^2 - 1)} \quad (7)$$

With:

$$D(b_i^2 \cdot L^2 - 1) \neq 0 \quad (8)$$

By introducing the coefficients A_α and B_α respectively in the Eq. (4), we obtain the different expressions of the excess minority carriers density.

And in pointing up the effect of the high recombination velocities to the junction, the profiles of the minority carriers' density in excess depending on the x depth in the base and according to the different illumination modes are represented in the following Fig. 2, 3 and 4.

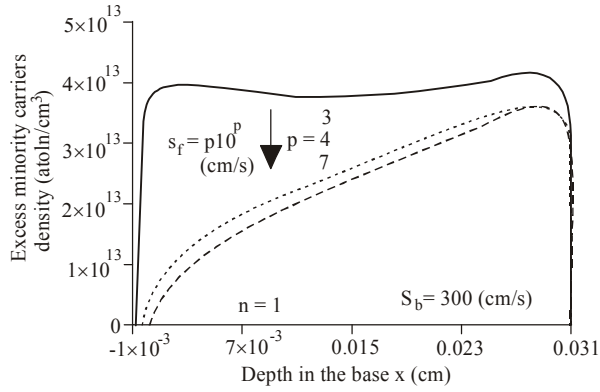


Fig. 4: Carriers' density according to the x depth in the base for both sides simultaneous illumination ($S_b = 3 \times 10^3$ cm/s, $D = 26$ cm²/s, $L = 0.015$ cm and $n = 1$)

According to the Eq. (5) which governs the boundary's conditions at the junction, the excess minority carriers density $\delta_a(x)$ is intimately linked at the recombination velocity of the excess minority carriers at the junction S_f and depth in the base. We have different characteristic regions when the solar cell is illuminated differently.

In Fig. 2, when the solar cell is illuminated in front side, the excess minority carriers density $\delta_1(x)$ increase and achieve a maxima point near the junction and decrease to minimum in the back side of the solar cell. When the solar cell is illuminated in rear side (Fig. 2), $\delta_2(x)$ increase from junction to a maxima point situated in the back side of the solar cell. We note that in Fig. 3, for the low values of junction recombination velocity S_f , the first maxima point is near the junction but for the high values of S_f , the second maxima point is at the rear side of the solar cell.

Concept of the recombination velocity to the junction initiating the short-circuit: From the expressions of the excess minority carriers density, in this part of the study, the survey deals with photocurrent densities profiles according to the recombination velocity to the junction among other parameters. These profiles are drawn on behalf of the three illuminations modes. From these profiles, we will study technics to determine the recombination velocity to the junction initiating the short circuit, its validity domain as well. And in a way these velocities expressed with a p parameter are put in the form of the following Eq. (9):

$$S_{f\alpha}(p) = p \cdot 10^p \quad (9)$$

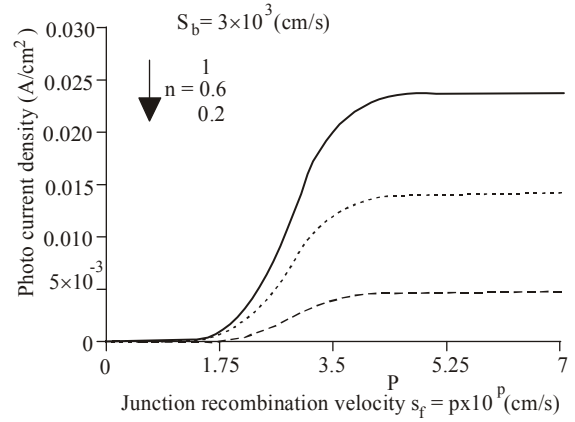


Fig. 5: Current density versus Log (S_f) for front side illumination and n illumination level ($S_{b1} = 3 \times 10^3$ cm/s, $D = 26$ cm²/s, $L = 0.015$ cm and $H = 0.03$ cm)

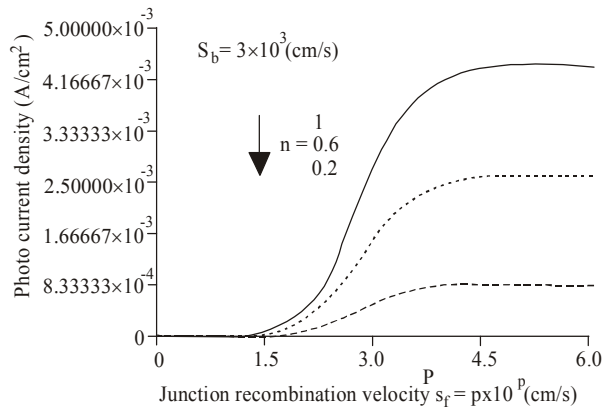


Fig. 6: Current density versus Log (S_f) for back side illumination and n illumination level ($S_{b1} = 3 \times 10^3$ cm/s, $D = 26$ cm²/s, $L = 0.015$ cm for three level of n)

Photocurrent density survey for the different light modes: The expression of the photocurrent density is obtained from the following relation:

$$J_{ph\alpha} = q \cdot D \cdot \left. \frac{\partial \delta_a(x)}{\partial x} \right|_{x=0} \quad (10)$$

where, $q = 1.6 \times 10^{19}$ C is the charge of the electron.

The expression (10) of the current density and on behalf of the three illumination modes ($\alpha = 1$ or 2 or 3) lets us draw its variation according to the recombination velocity at the junction and for different values of the illumination level.

We plot in Fig. 5, 6 and 7 photocurrent density versus Log (S_f) respectively for front side, back side and simultaneous n illumination level.

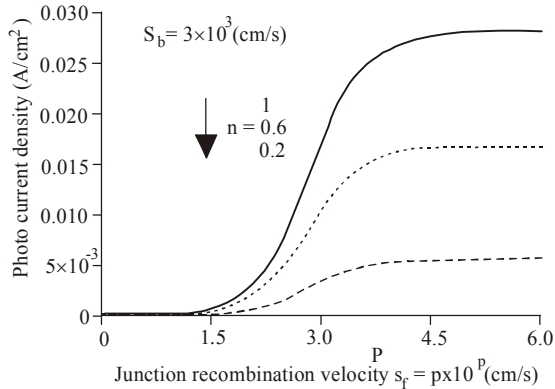


Fig. 7: Current density versus Log (S_f) for simultaneous illumination and n illumination level ($S_{b1} = 3 \times 10^3$ cm/s, $D = 26$ cm²/s, $L = 0.015$ cm and $H = 0.03$ cm, for three level of n)

The curves present two stages for all three figures: We note that the photocurrent's density increases with the recombination velocity up to a value where it remains constant whatever the illuminations mode and level. This constant value of photocurrent is called the short-circuit photocurrent. It depends on illumination mode and occurs when the junction recombination velocity is greater than 4×10^4 cm/s as shown in Diallo *et al.* (2008).

For the low values of S_f which are smallest than $1.5 \times 10^{1.5}$ cm/s, the photocurrent density is near zero.

Recombination velocity to the junction S_{fcca} for different modes and levels of the solar cell's light:

The condition of short-circuit is a situation of total shift to the emitter of all the electrons photo-generated in the base and that cross the junction. And according to the currents' profile below, this situation of short-circuit appears to the high values of the recombination velocity to the junction S_{fa} . For such a situation, there is not any longer some carriers to the junction storing and the photocurrent is maximal and tends asymptotically to a constant value.

This fact is mathematically due to the following equation:

$$\frac{\partial J_{pha}}{\partial S_{fa}} = 0 \tag{11}$$

In other terms, the general expression of the short-circuit's current's density is (Ba *et al.*, 2003):

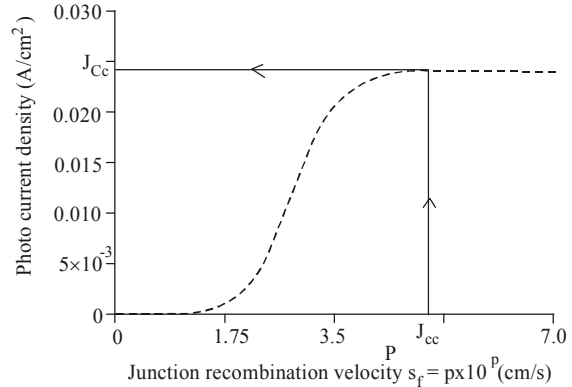


Fig. 8: Technics to determine the short-circuit current

$$J_{cca} = \lim_{S_{fa} \rightarrow \infty} J_{pha} \tag{12}$$

The first intercept point between the calibrated function $J_{pha}(S_{fa})$ and the short circuit current J_{cca} yields the value of the recombination velocity to the junction. So, this will be deduced from the following equation for the three illumination modes:

$$J_{pha}(S_{fa}) - J_{cca} = 0 \tag{13}$$

Technics to determine short-circuit current: The Fig. 8 elucidates this technics that lets determine the short-circuit current for a recombination velocity to the given junction. We focus on a particular velocity of recombination to the junction. It's the one that initiates the short-circuit current and in the example that we have given, it is equal to $S_{fcc} = 4.9 \times 10^{4.9}$ cm/s. The projection of the line $X = p = 4.9 \times 10^{4.9}$ on the current variation curve $J_{ph}(n, p, m)$ according to the recombination velocity S_f we give an intersection point. The projection of this intersection point Y-axis gives us the precise value of the short-circuit current.

However, this technics to determine the short-circuit current has some limits. Particularly, for some recombination velocities to the junction inferior to the ones initiating the short-circuit current, this technics is not usable. Its validity domain is defined for recombination velocities superior or equal to a reference velocity and in our case $4.9 \times 10^{4.9}$ cm/s.

Recombination velocity to the junction: In resolving the Eq. (13) that becomes (14), we can deduce from there the recombination velocity at the junction for the three illumination modes.

Table 2: Recombination velocity to the junction S_{fcc1} initiating the short-circuit current for solar cell's front side illuminated

Light level (n)	1	0.6	0.2
Short circuit current J_{cc1} (A)	$23.53.10^{-3}$	$14.11.10^{-3}$	$4.67.10^{-3}$
Velocity S_{b1} (cm / s)	3.10^3	3.10^3	3.10^3
S_{fcc1} (cm / s) initiating the short circuit current J_{cc1} (A)	$3.89.10^6$	$2.09.10^6$	$5.46.10^5$

Table 3: Recombination velocity to the junction S_{fcc2} initiating the short-circuit current for a back side illumination

Light level (n)	1	0.6	0.2
Short circuit current J_{cc2} (A)	$4.30.10^{-3}$	$2.57.10^{-3}$	$8.55.10^{-4}$
Velocity S_{b2} (cm/s)	3.10^3	3.10^3	3.10^3
S_{fcc2} (cm/s) initiating the short circuit current J_{cc2} (A)	$2.30.10^5$	$1.52.10^5$	$1.33.10^5$

Table 4: The recombination velocity to the junction S_{fcc3} initiating the short-circuit current for a simultaneous light

Light level (n)	1	0.6	0.2
Short circuit current J_{cc3} (A)	$27.96.10^{-3}$	$16.77.10^{-3}$	$5.58.10^{-3}$
Velocity S_{b3} (cm/s)	3.10^3	3.10^3	3.10^3
S_{fcc3} (cm/s) initiating the short circuit current J_{cc3} (A)	$4.23.10^6$	$2.26.10^6$	$6.76.10^5$

$$J_{p\alpha}(S_f) - J_{cca} = qS_f \left(\frac{B_\alpha}{L} - \sum_{i=1}^3 K_i \cdot b_i \left(-\xi_\alpha + \chi_\alpha e^{-b_i H} \right) \right) - J_{cd} = 0 \quad (14)$$

In these conditions, Eq. (14) lets us obtain the junction recombination velocity's expression which depends to the illumination level n, the highest value of the short-circuit current J_{cca} and the back side recombination velocity S_{ba} .

So, the recombination velocity at the junction is expressed as follows:

$$S_{fcc\alpha}(n, J_{cca}, m) = \frac{\psi(n, J_{cca}) \cdot D^2 \cdot \sinh\left(\frac{H}{L}\right) + M}{LD \cdot X(m) \sum_i^3 K_i - L \sum_{i=1}^3 K_i \cdot E_i(i, m) - F} \quad (15)$$

- **For a front side light:** So, we can establish the Table 2 linking some parameters of the solar cell when the latter is illuminated by its emitter. According to this table, for a front side illumination, we note that the recombination

velocity to the junction as the short-circuit current grows with the light level. This result confirms once the proportionality between the charge carrier density and the light level. We also note that the short circuit current grows with the recombination velocity to the junction.

- **Recombination velocity to the junction S_{fcc2} for a back side light:** Some parameters of the solar cell are given in Table 3, for back side illumination. By this Table 3 for a back side light, the analysis on the short-circuit current like on the recombination velocity to the junction remains the same as the solar cell's light by the emitter. Still, we note that these values are less important than for the front side light.
- **Recombination velocity to the junction for a simultaneous light:** The Table 4 links some parameters of the solar cell for its simultaneous light.

The interpretation of the Table 4 is identical for the front and back side illuminations. However, we note that the short-circuit current like the recombination velocities to the junction for a simultaneous illumination are more important than for the two other modes.

CONCLUSION

This study has made it possible for us to show the dependence of the photocurrent to the illumination level. No matter the number of sun, we note that the photocurrent density increases with the recombination velocity at the junction up to some value where it remains constant. Then the short-circuit condition corresponds to some S_{fcc} high values where the photocurrent tends asymptotically towards a constant value giving the short-circuit current. This situation of short-circuit corresponding to some recombination velocities very high to the junction transforms the solar cell into a current generator.

In this study we have also proposed a technic to determine the first value S_{fcc} of the recombination velocity yielding the short-circuit current and its validity domain.

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